

Magnetic Induction Eq. (cont.)

if we ignore Hall term & Polarization E field, we have

$$\vec{E} = -\frac{1}{c} \vec{v} \times \vec{B} + \frac{1}{\sigma} \vec{J}$$

then, Faraday's law gives

$$\frac{\partial \vec{B}}{\partial t} = -c \nabla \times \vec{E}$$

$$= -c \nabla \times \left(-\frac{1}{c} \vec{v} \times \vec{B} + \frac{1}{\sigma} \vec{J} \right)$$

$$= \nabla \times (\vec{v} \times \vec{B}) - \frac{c}{\sigma} \nabla \times \vec{J}$$

} Ampere's law

$$= \nabla \times (\vec{v} \times \vec{B}) - \frac{c^2}{4\pi\sigma} \nabla \times (\nabla \times \vec{B})$$

} vec. ident

$$= \nabla \times (\vec{v} \times \vec{B}) + \frac{c^2}{4\pi\sigma} \nabla^2 \vec{B}$$

advection term

diffusion term

define Magnetic Reynolds #

R_M

$$= \frac{\text{advection term}}{\text{diffusion term}} = \frac{\frac{1}{L} U B}{\frac{c^2}{4\pi\sigma} \frac{B}{L^2}}$$

$$= \frac{LU 4\pi\sigma}{c^2} = \frac{LU}{\eta}$$

$$\eta = \frac{c^2}{4\pi\sigma} \text{ mag. diffus.}$$

What is R_m in this room?

$$R_m = \frac{LU4\pi\sigma}{c^2}$$

$$L \approx 1 \text{ m} \approx 100 \text{ cm}$$

$$U \approx 10 \text{ mph} \approx 450 \frac{\text{cm}}{\text{s}}$$

$$\sigma_{\text{air}} \approx 4.5 \times 10^{-5} \frac{1}{\text{s}} \text{ (c.g.s.)}$$

$$= \frac{(100)(450)(4\pi)(4.5 \times 10^{-5})}{(3 \times 10^{10})^2}$$

$$= 3 \times 10^{-20} \ll 1$$

note also, the plasma β in this room is

$$\beta = \frac{P}{(B^2/8\pi)} = \frac{1 \text{ bar}}{B^2/8\pi} = \frac{10^6 \frac{\text{dyne}}{\text{cm}^2}}{\frac{(1)^2}{8\pi} \frac{\text{dyne}}{\text{cm}^2}}$$

$$\approx 8\pi \times 10^6 \gg 1$$

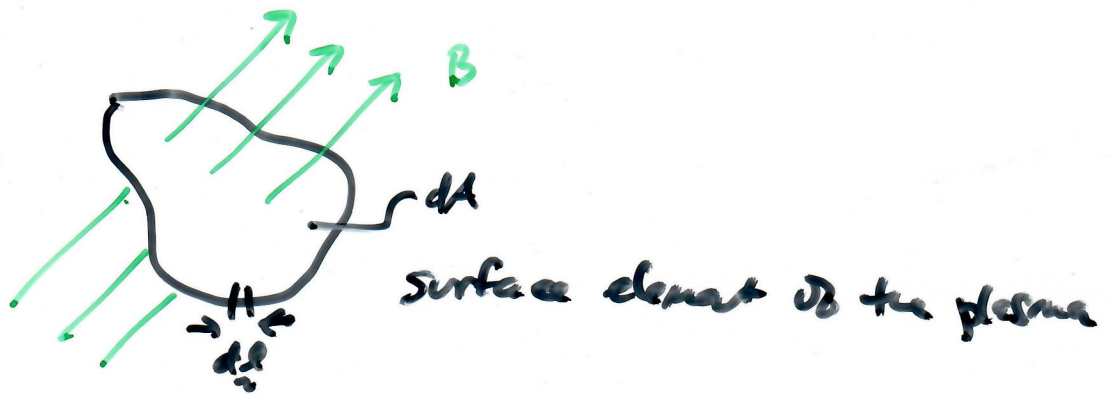
But, in ^{most} Astrophysical / Heliospheric plasmas

$$R_m \gg 1$$

Frozen-flux theorem

$$\frac{\partial \underline{B}}{\partial t} = \nabla \times (\underline{u} \times \underline{B}) + \eta \nabla^2 \underline{B}$$

if $\eta \rightarrow 0$, \underline{B} is "frozen" into the plasma flow



flux through area is $\int \underline{B} \cdot d\underline{A}$

$$\begin{aligned} \frac{\partial}{\partial t} \left(\int \underline{B} \cdot d\underline{A} \right) &= \int \frac{\partial \underline{B}}{\partial t} \cdot d\underline{A} + \int \underline{B} \cdot \frac{d\underline{A}}{dt} \quad \left. \begin{array}{l} \text{Leibnitz} \\ \text{theorem} \end{array} \right\} \\ &= \int \frac{\partial \underline{B}}{\partial t} \cdot d\underline{A} + \oint \underline{B} \cdot (\underline{u} \times d\underline{l}) \\ &= \int \frac{\partial \underline{B}}{\partial t} \cdot d\underline{A} + \oint (\underline{B} \times \underline{u}) \cdot d\underline{l} \quad \left. \begin{array}{l} \text{vec} \\ \text{identity} \end{array} \right\} \\ &= \int \frac{\partial \underline{B}}{\partial t} \cdot d\underline{A} - \oint (\underline{u} \times \underline{B}) \cdot d\underline{l} \end{aligned}$$

4-
strong
f.m.

$$= \int \frac{\partial \underline{B}}{\partial t} \cdot d\underline{A} - \int [\nabla \times (\underline{u} \times \underline{B})] \cdot d\underline{A}$$

$$= \int \left[\frac{\partial \underline{B}}{\partial t} - \nabla \times (\underline{u} \times \underline{B}) \right] \cdot d\underline{A}$$

$$= 0 \quad \text{if} \quad \eta = 0$$

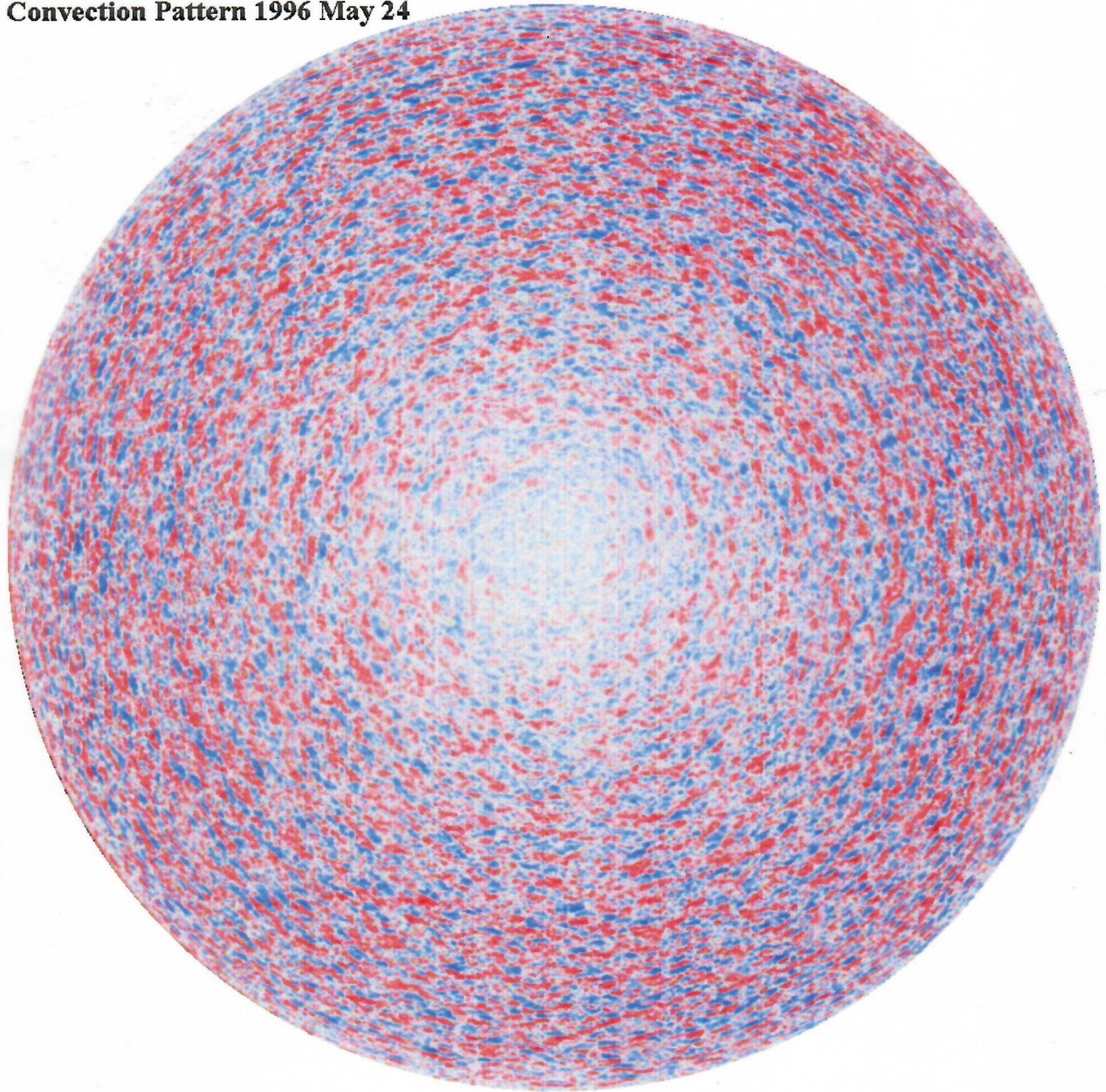
the $\nabla \times (\underline{u} \times \underline{B})$ term in the mag. induction eq. leads to field amplification

- turbulent dynamos \rightarrow strong field creation
- α - ω dynamos \rightarrow generate stellar & planetary fields

an example of a turbulent dynamo is the sun's "network" magnetic field

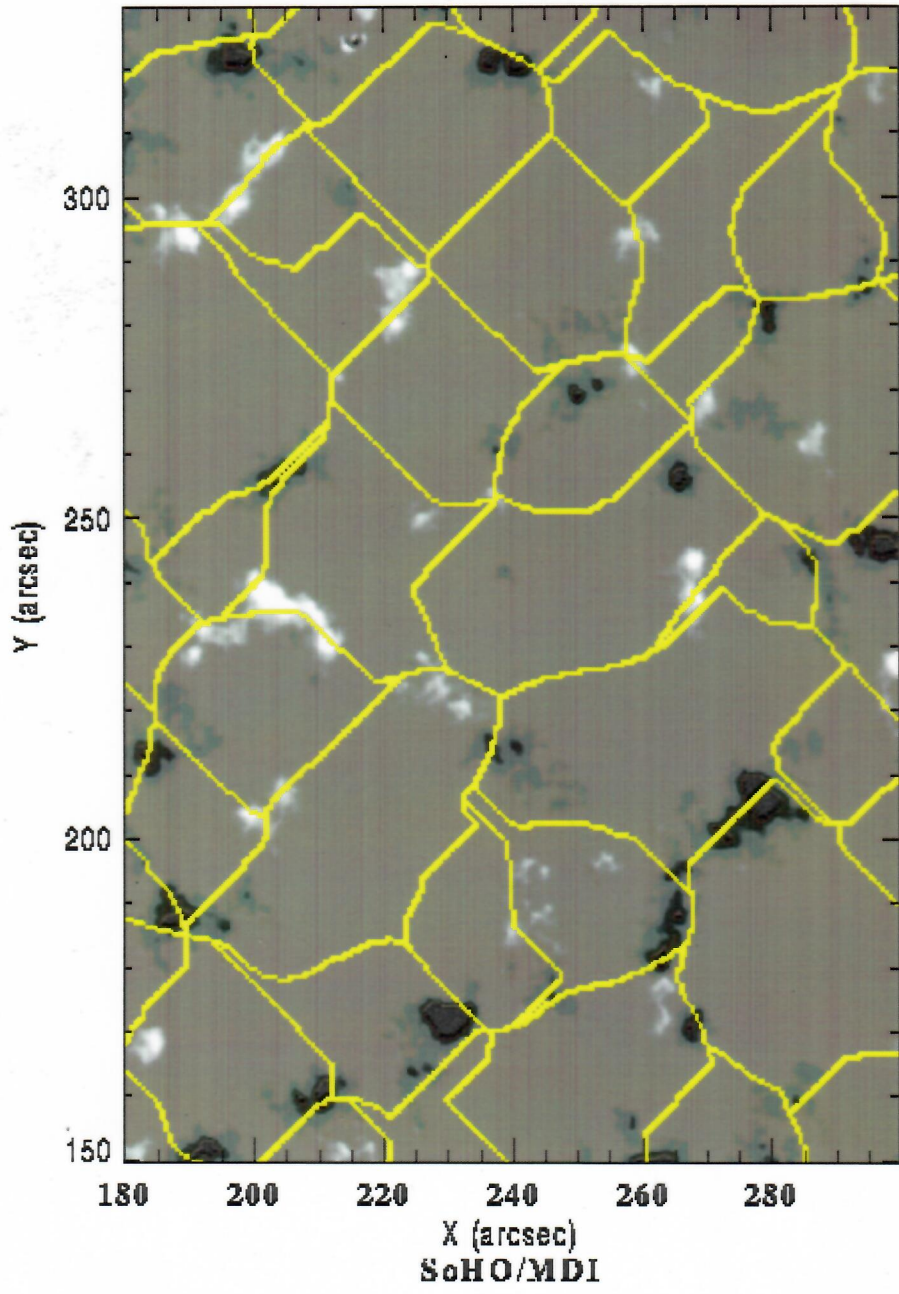
6-10
5-10
10-10

Convection Pattern 1996 May 24



~~2~~ ~~20~~
~~6~~ ~~20~~

23 Feb. 1996, 16:44 to 21:03 UT

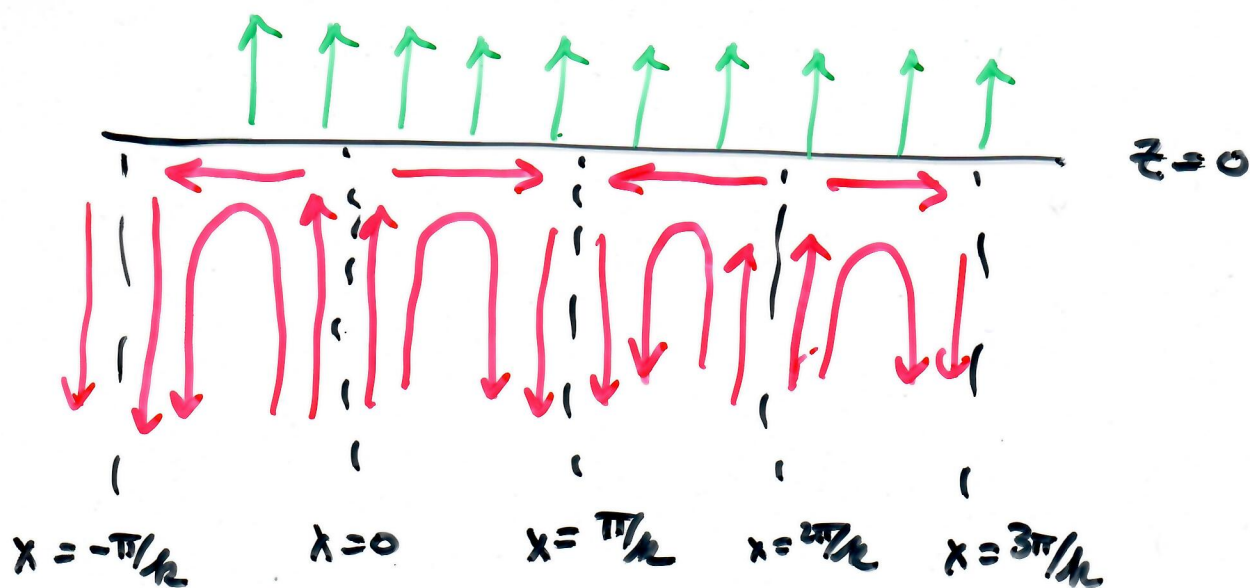


Consider the following kinematic representation of the transverse plasma flow at the Sun due to supergranulation

$$\vec{u} = (u_x, 0, u_z) \quad \vec{B} = (0, 0, B(x))$$

$$u_x = u_0 \sin kx$$

$$u_z = -kz \cos kx$$



Initially $\frac{B^2}{8\pi} \ll \rho u^2$ $\rho = \text{plasma density}$

How does $B(x,t)$ evolve with x & t ?

$$\frac{\partial \underline{B}}{\partial t} = \nabla \times (\underline{u} \times \underline{B})$$

ignore mag. diffusi:

$$\underline{B} = B(x, t) \hat{z}$$

u_x, B do not depend on z

$$\frac{\partial \underline{B}}{\partial t} = \frac{\partial}{\partial z} (u_x B) \hat{x} - \frac{\partial}{\partial x} (u_x B) \hat{z}$$

$$\Rightarrow \frac{\partial B(x, t)}{\partial t} = -\frac{\partial}{\partial x} (u_x B)$$